

Lecture 3: Microscopic models of the dispersion

Petr Kužel

- Origins of the dispersion
 - Dielectric resonances: bounded (polarization) charges
 - Elastically bounded electrons (exciton polaritons)
 - Polar lattice vibrations (phonon polaritons)

$$m\ddot{\mathbf{r}} = -m\gamma\dot{\mathbf{r}} - k\mathbf{r} - e\mathbf{E}$$

- Metallic resonances: free charges
 - Conductivity in metals (underlying lattice resonances are much weaker)

$$m\ddot{\mathbf{r}} = -m\gamma\dot{\mathbf{r}} - e\mathbf{E}$$

Electromagnetic force on charges

$$f = -e(\mathbf{E} + \mathbf{v} \wedge \mathbf{B}) = -e\left(\mathbf{E} + \frac{\mathbf{v}}{c} \wedge (s \wedge \mathbf{E})\right) \approx -e\mathbf{E}$$

- Only electric part in the non-relativistic approximation

Dielectric resonance

$$\ddot{\mathbf{P}} + \gamma\dot{\mathbf{P}} + \frac{k}{m}\mathbf{P} = \frac{Ne^2}{m}\mathbf{E}_{loc} \quad (\mathbf{P} = -Ner), \quad \mathbf{E}_{loc} = \frac{\mathbf{P}}{3\epsilon_0} + \mathbf{E}$$

$$\ddot{\mathbf{P}} + \gamma\dot{\mathbf{P}} + \omega_0^2\mathbf{P} = \frac{Ne^2}{m}\mathbf{E} \quad \omega_0 = \sqrt{\frac{k}{m} - \frac{Ne^2}{3m\epsilon_0}}$$

Electric field is decomposed to harmonic waves:

$$\mathbf{P} = \frac{Ne^2/m}{\omega_0^2 - \omega^2 - i\omega\gamma}\mathbf{E} \quad \mathbf{D} = \epsilon_0\mathbf{E} + \mathbf{P} = \epsilon_0\left(1 + \frac{Ne^2}{m\epsilon_0} \frac{1}{\omega_0^2 - \omega^2 + i\omega\gamma}\right)\mathbf{E} = \epsilon\mathbf{E}$$

Dispersion near dielectric resonances

$$\epsilon_r = 1 + \frac{Ne^2}{m\epsilon_0} \frac{1}{\omega_0^2 - \omega^2 + i\omega\gamma} = 1 + \frac{\omega_L^2 - \omega_0^2}{\omega_0^2 - \omega^2 + i\omega\gamma}, \quad \epsilon_{stat} = \frac{\omega_L^2}{\omega_0^2}$$

Small damping ($\omega_0 \gg \gamma$):

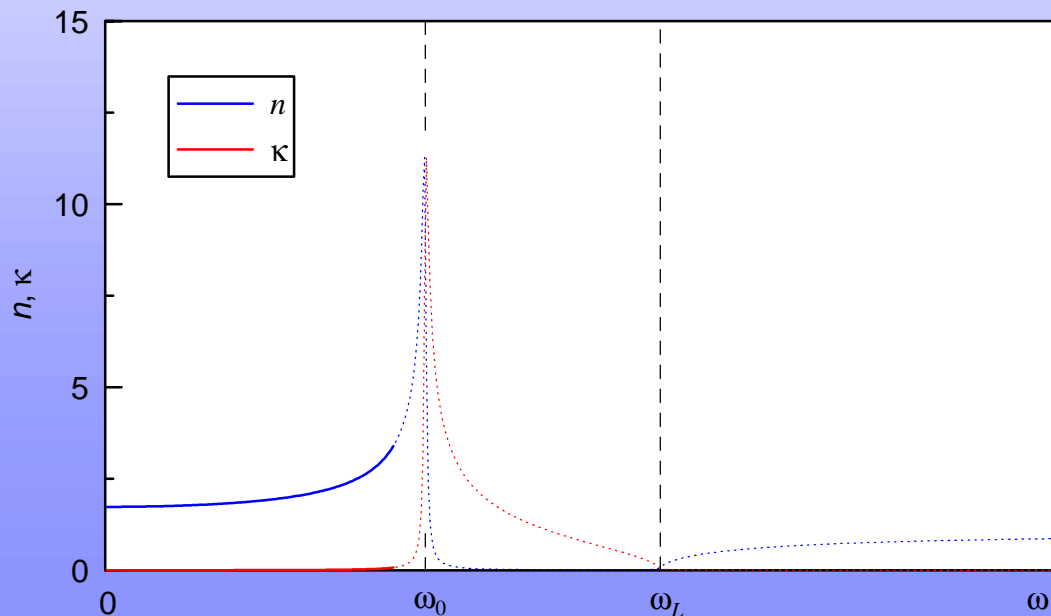
- 1) Low-frequency regime ($0 \leq \omega \leq \omega_0 - 2\gamma$)
- 2) Vicinity of the resonance ($\omega_0 - 2\gamma \leq \omega \leq \omega_0 + 2\gamma$)
- 3) High reflectivity regime (Reststrahlung band)
($\omega_0 + 2\gamma \leq \omega \leq \omega_L$)
- 4) High-frequency regime ($\omega \geq \omega_L$)

Dispersion: dielectric resonances

- 1) Low-frequency regime ($0 \leq \omega \leq \omega_0 - 2\gamma$), $n^2 = \epsilon_r$, $\kappa = 0$, normal dispersion

$$\epsilon_r = 1 + \frac{\omega_L^2 - \omega_0^2}{\omega_0^2 - \omega^2 + i\omega\gamma} \approx \epsilon_{stat} + \frac{\omega^2}{\omega_0^2} (\epsilon_{stat} - 1)$$

$$\begin{aligned} \epsilon_r' &= n^2 - \kappa^2 \\ \epsilon_r'' &= 2n\kappa \end{aligned}$$

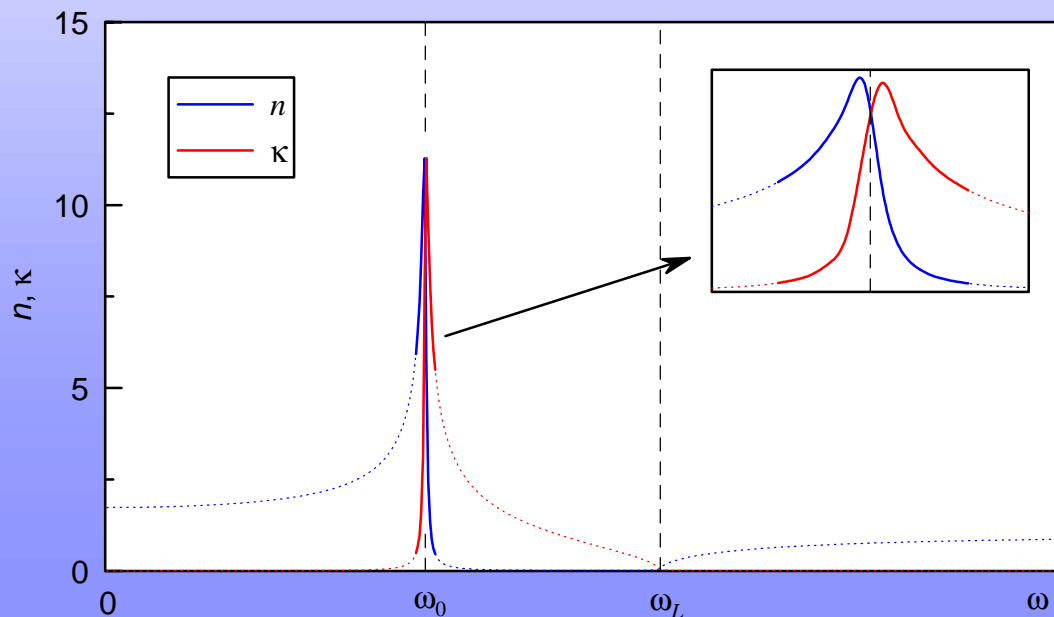


Dispersion: dielectric resonances

- 2) Vicinity of the resonance ($\omega_0 - 2\gamma \leq \omega \leq \omega_0 + 2\gamma$),
anomalous dispersion

$$\epsilon_r = 1 + \frac{\omega_L^2 - \omega_0^2}{\omega_0^2 - \omega^2 + i\omega\gamma} \approx 1 - \frac{\omega_L^2 - \omega_0^2}{\omega_0} \frac{2\Delta + i\gamma}{4\Delta^2 + \gamma^2}$$

$$\begin{aligned}\epsilon_r' &= n^2 - \kappa^2 \\ \epsilon_r'' &= 2n\kappa\end{aligned}$$

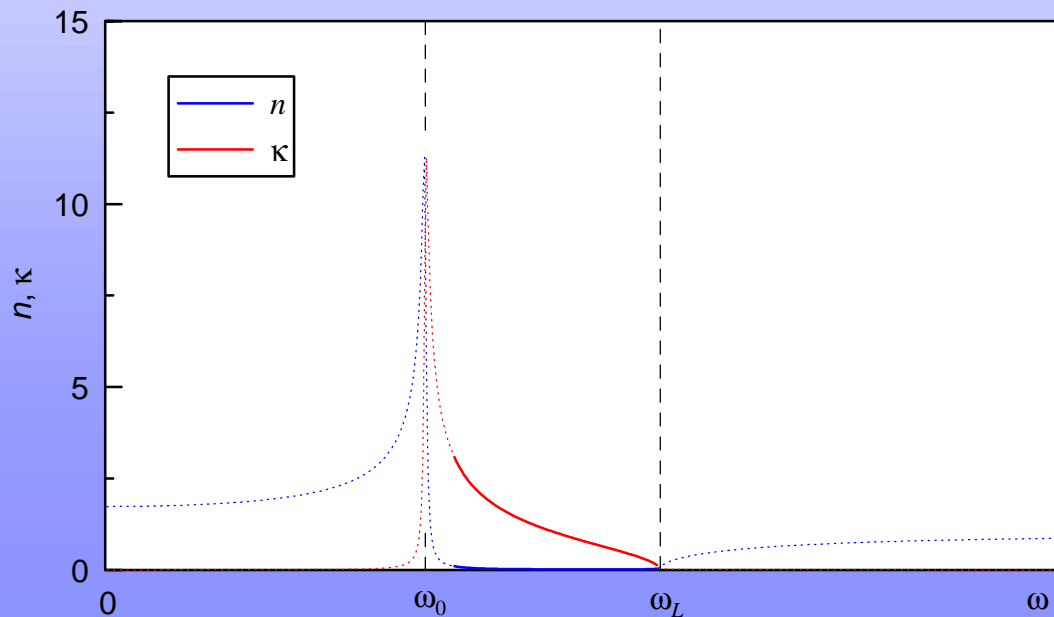


Dispersion: dielectric resonances

3) High reflectivity regime ($\omega_0 + 2\gamma \leq \omega \leq \omega_L$), $\kappa^2 = \epsilon_r$, $n = 0$

$$\epsilon_r = 1 + \frac{\omega_L^2 - \omega_0^2}{\omega_0^2 - \omega^2 + i\omega\gamma} \approx 1 - \frac{\omega_L^2 - \omega_0^2}{\omega^2 - \omega_0^2} < 0$$

$$\begin{aligned} \epsilon_r' &= n^2 - \kappa^2 \\ \epsilon_r'' &= 2n\kappa \end{aligned}$$

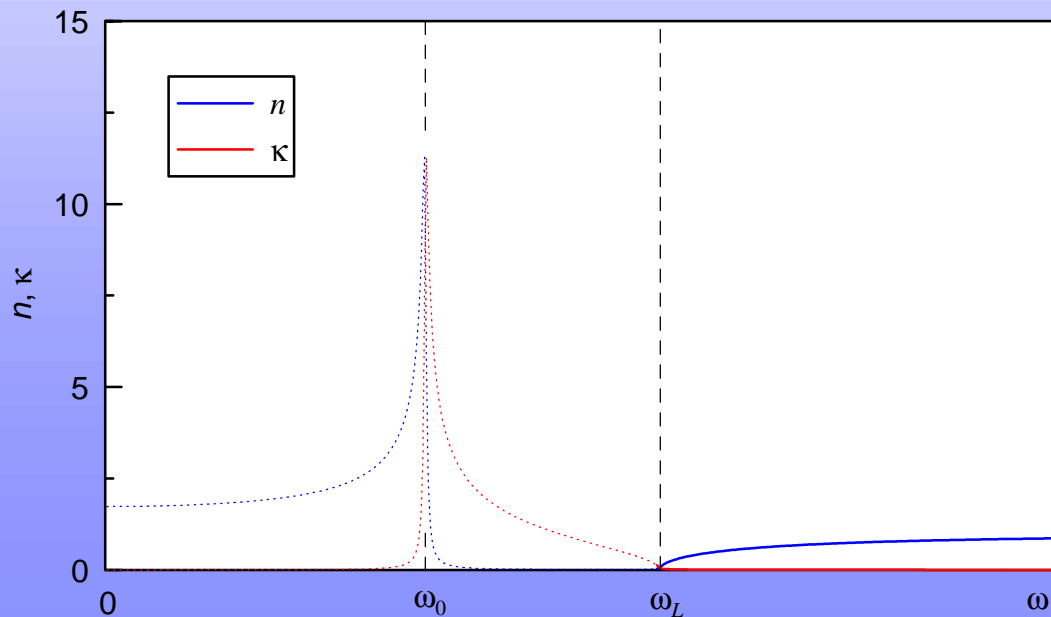


Dispersion: dielectric resonances

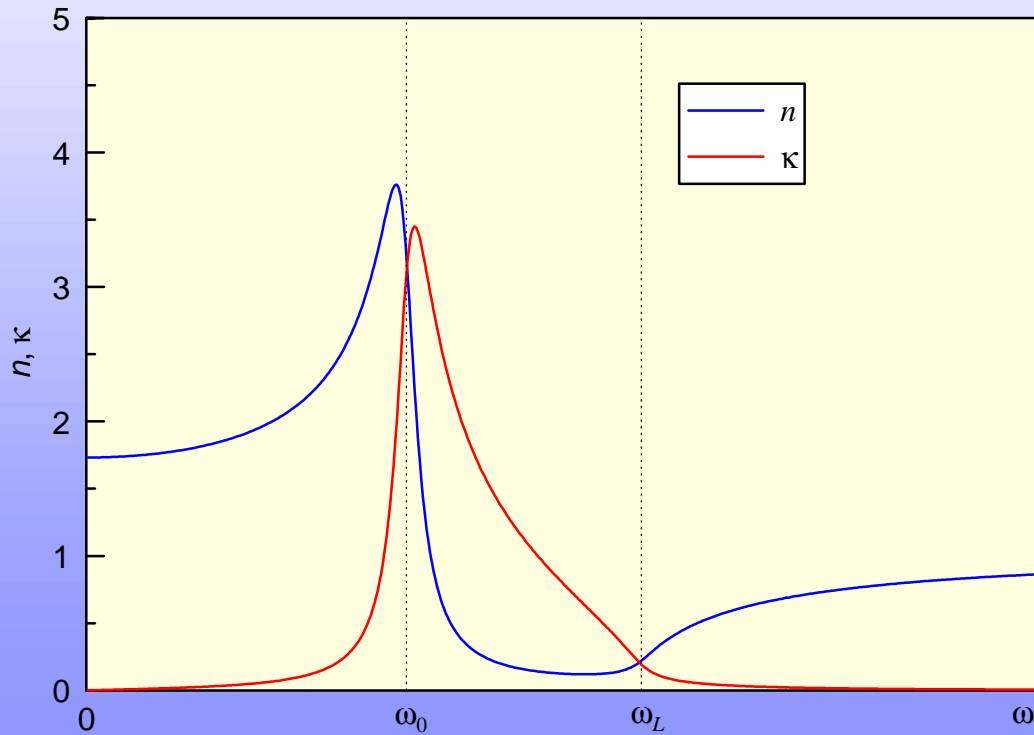
4) High-frequency regime ($\omega \geq \omega_L$), normal dispersion

$$\epsilon_r = 1 + \frac{\omega_L^2 - \omega_0^2}{\omega_0^2 - \omega^2 + i\omega\gamma} \approx 1 - \frac{\omega_L^2 - \omega_0^2}{\omega^2 - \omega_0^2} > 0$$

$$\begin{aligned}\epsilon_r' &= n^2 - \kappa^2 \\ \epsilon_r'' &= 2n\kappa\end{aligned}$$

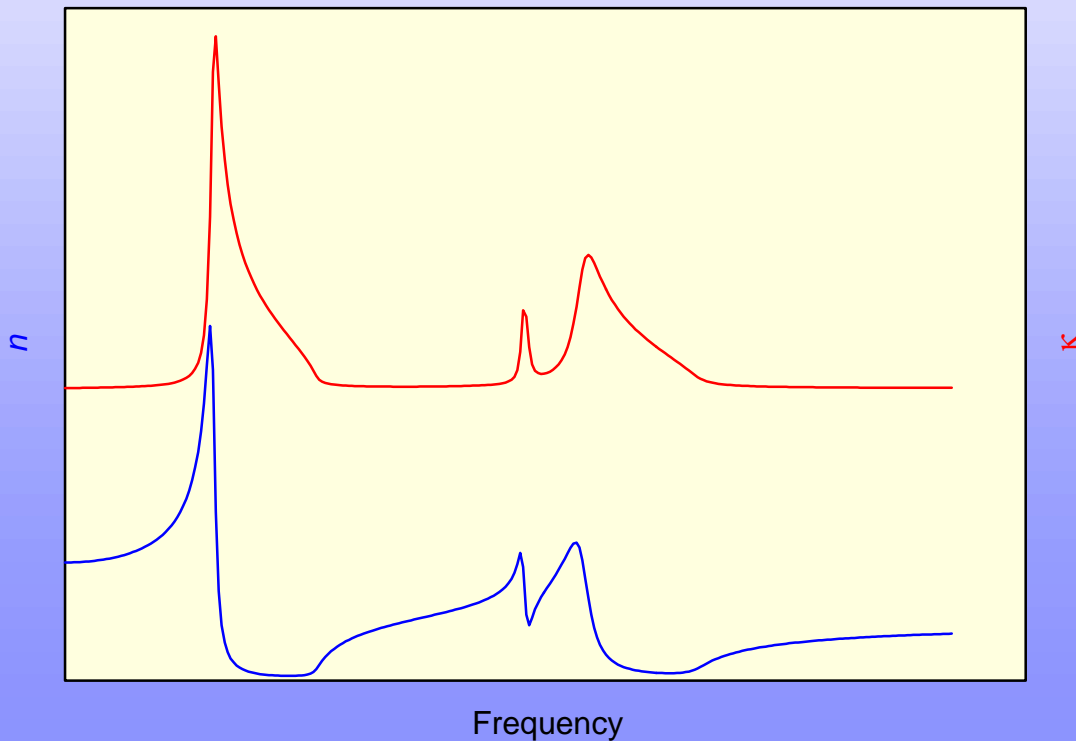


Dielectric resonance: stronger damping



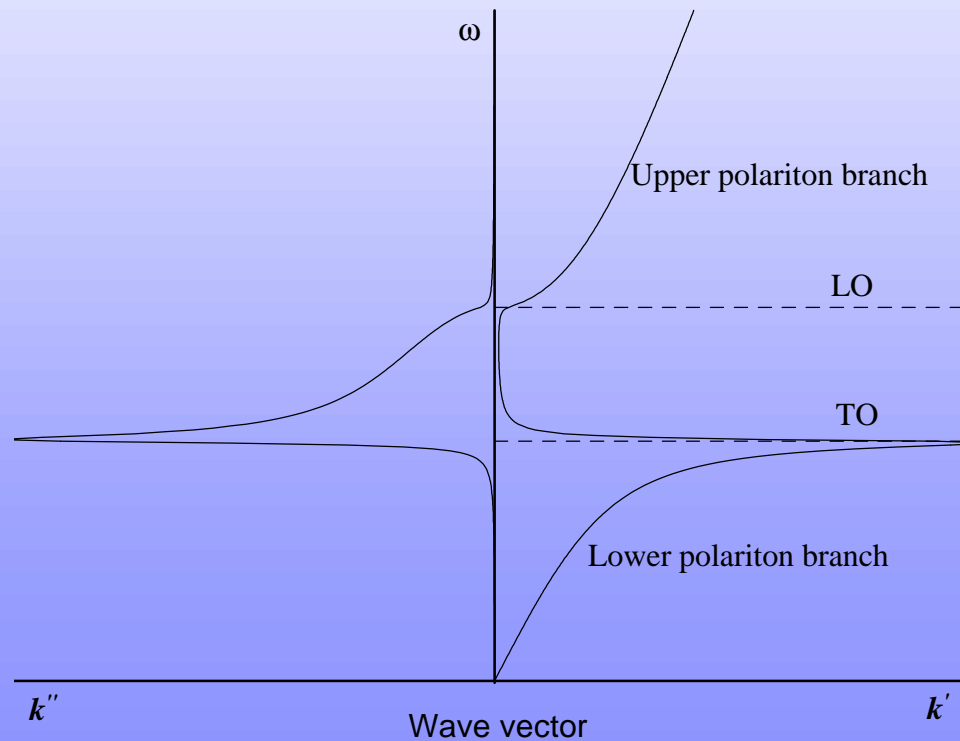
Several dielectric resonances

$$\epsilon = \epsilon_0 \left(1 + \sum_j \frac{f_j}{\omega_j^2 - \omega^2 + i\omega\gamma_j} \right)$$



Polariton concept

Polariton equation: $k = \frac{\omega}{c} n(\omega)$



Metallic resonance

$$\dot{\mathbf{P}} + \gamma \mathbf{P} + \frac{k}{m} \mathbf{P} = \frac{Ne^2}{m} \mathbf{E}_{loc} \quad (\mathbf{P} = -Ner), \quad \mathbf{E}_{loc} = \frac{\mathbf{P}}{3\epsilon_0} + \mathbf{E}$$

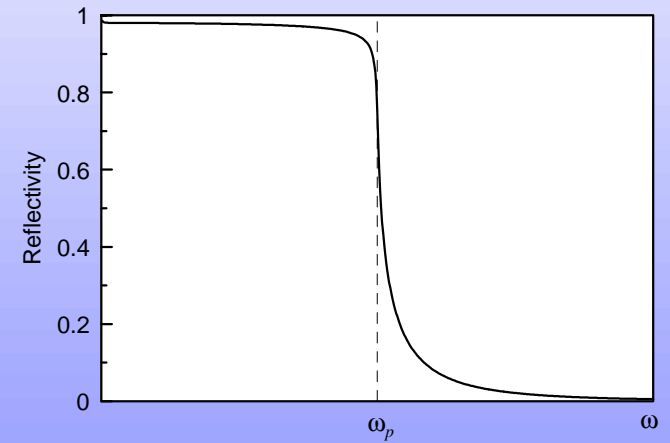
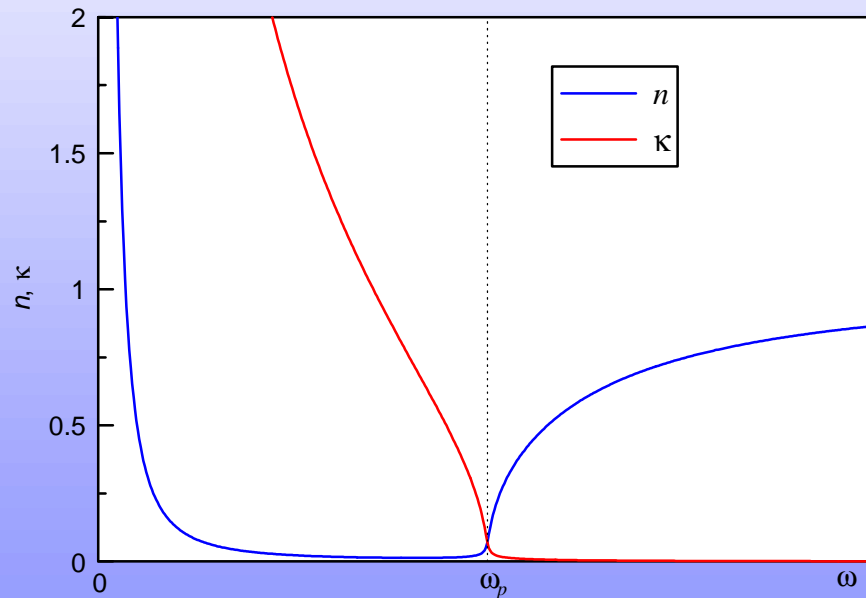
$$\frac{d\mathbf{j}}{dt} + \gamma \mathbf{j} = \left(\frac{Ne^2}{m} \right) \mathbf{E}$$

$$\sigma = \frac{Ne^2/m}{\gamma + i\omega} = \frac{\sigma_0}{1 + i\omega/\gamma}$$

$$\epsilon = \epsilon_0 \left(1 + \frac{Ne^2}{m\epsilon_0} \frac{1}{i\omega\gamma - \omega^2} \right) = \epsilon_0 \left(1 + \frac{\omega_p^2}{i\omega\gamma - \omega^2} \right)$$

ω_p ... plasma frequency (longitudinal resonance of the free charges)

Dispersion: metallic resonance



Metallic mirrors

